

Accelerator Physics Lecture Modules

- 1. Introduction/Timing**
- 2. Beam propagation basics**
- 3. Simple Harmonic Oscillator**
- 4. Betatron oscillations (transverse motion)**
- 5. Dispersion**
- 6. Closed Orbit Distortion**
- 7. Synchrotron oscillations**
- 8. Properties of synchrotron radiation**
- 9. Radiation damping**
- 10. Equilibrium emittance**

Synchrotron Radiation - The Dipole Oscillator

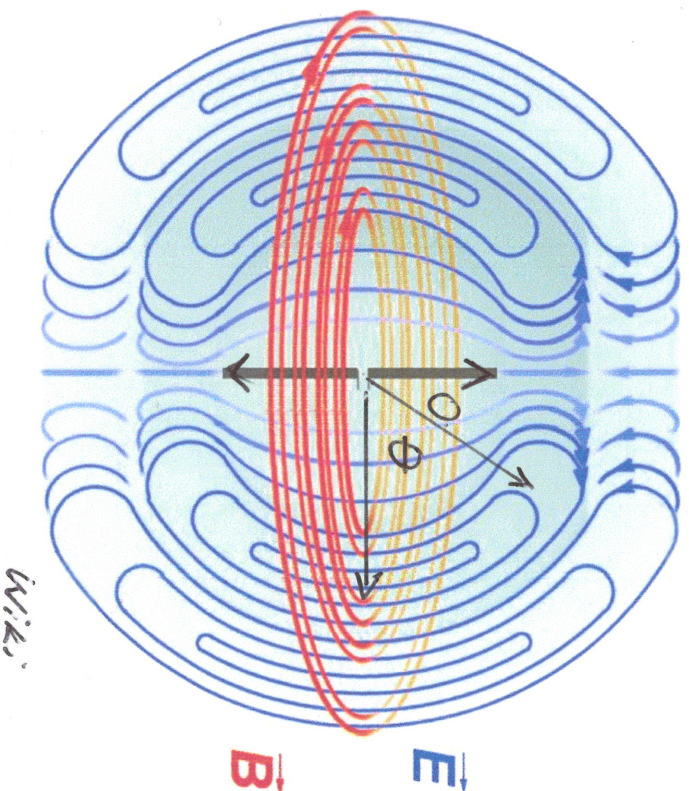
Examine field emission from a dipole (Radio Antenna)

- Charge acceleration is vertical
- Propagating field components transverse
- Poynting's Vector $S = \frac{c}{4\pi} [\vec{E} \times \vec{B}] \text{ W/m}^2$
- $E_L = \frac{q \cdot a}{c^2 R} \cdot \cos\theta$

$$S = \frac{q^2 a^2}{4\pi c^3 R^2} \cos^2\theta \quad \text{W/m}^2$$

For synchrotron radiation

$$F = \gamma m a = \frac{q \vec{v} \times \vec{B}}{c} \rightarrow a = \frac{q c B}{\gamma m c}$$



Dipole Oscillator Power

3-D field pattern looks like a doughnut $\rightarrow$

SR emission in particle system

$$S \propto \frac{q^2 a^2}{R^2} \cos^2 \theta \text{ W/m}^2$$

$$a = \frac{qB}{\gamma m c} \leftarrow \begin{array}{l} \text{strong field} \\ \text{light particle} \end{array}$$

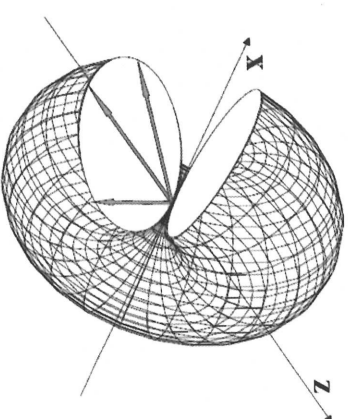
Integrate over doughnut surface
to get Watts (Joules/sec)

$$P_{rad} \propto E^2 B^2 \quad \text{or} \quad P_{rad} \propto \frac{E^4}{R^2}$$

$\nearrow$ particle energy $\nwarrow$ magnetic field

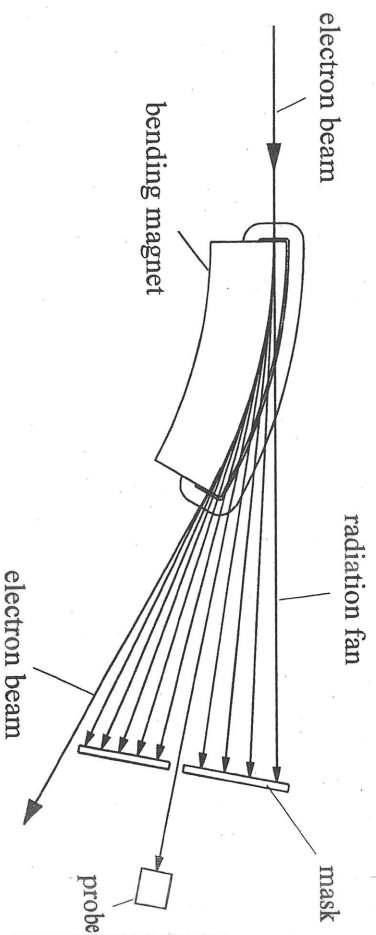
For a given bend radius

$$P_{rad} \propto E^4$$

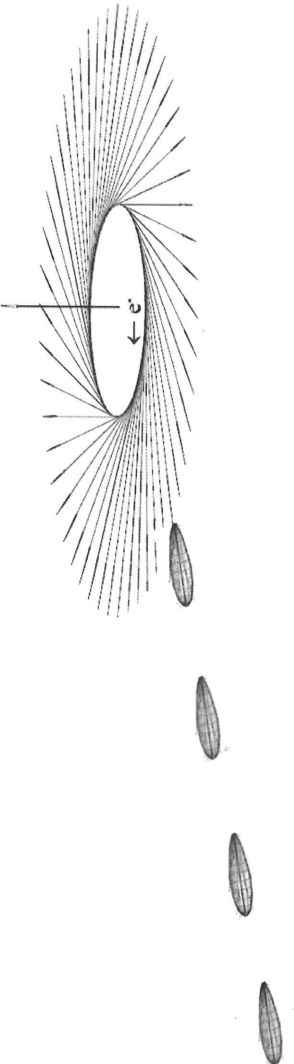


Synchrotron Radiation in a Storage Ring

Dipole bending magnets supply transverse acceleration

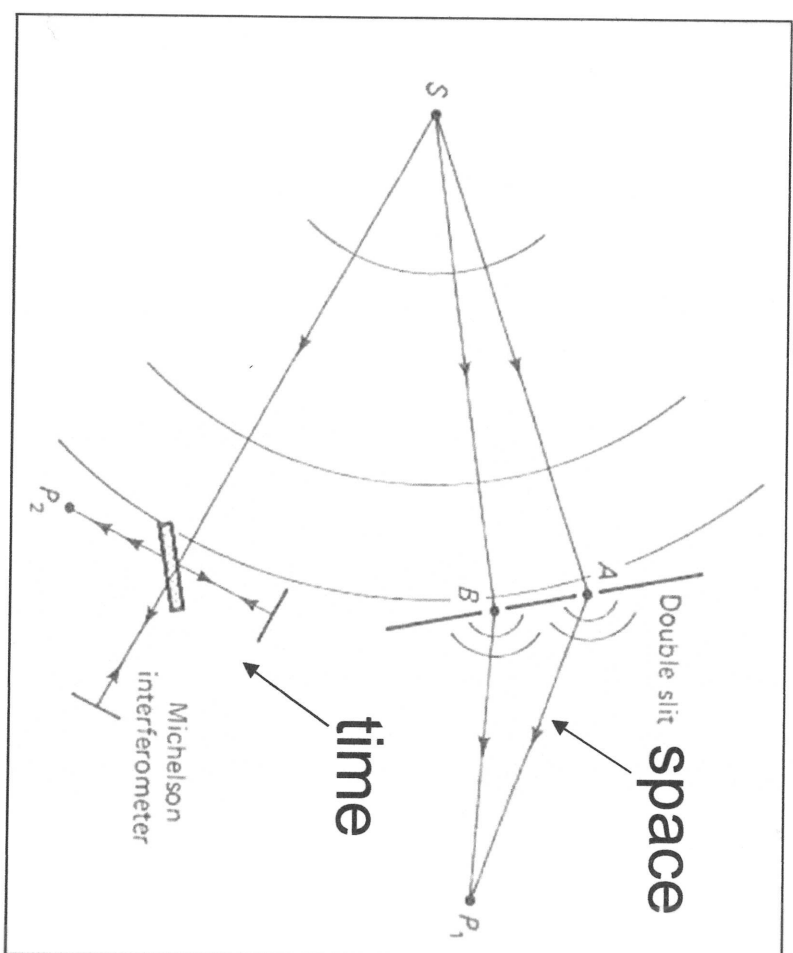


SR emission over 2π circumference with pulse structure



Note: individual photon emission events are quantized like an atom
radiation is longitudinally incoherent

Space and Time Correlation (Coherence)



At the beach...

SR Emission in the Laboratory Frame

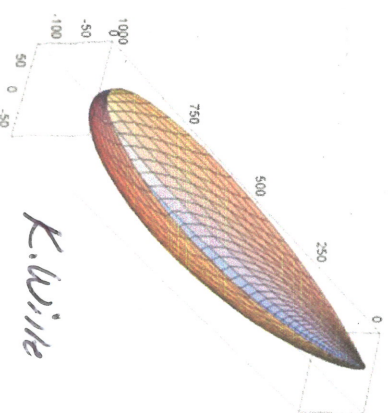
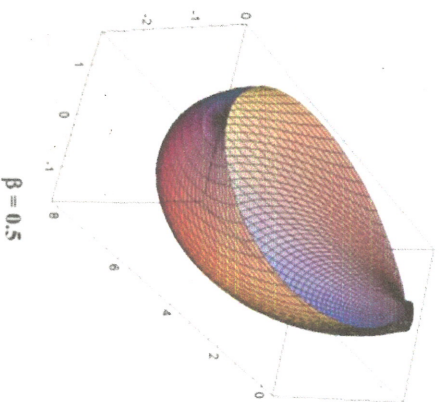
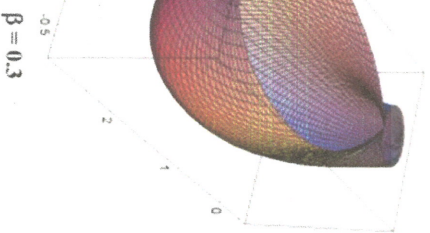
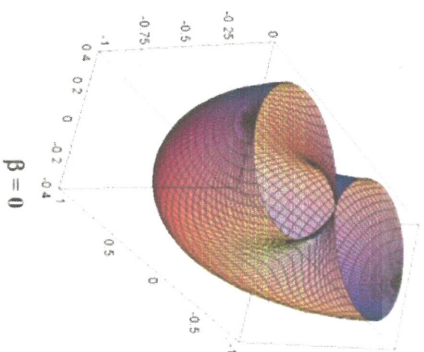
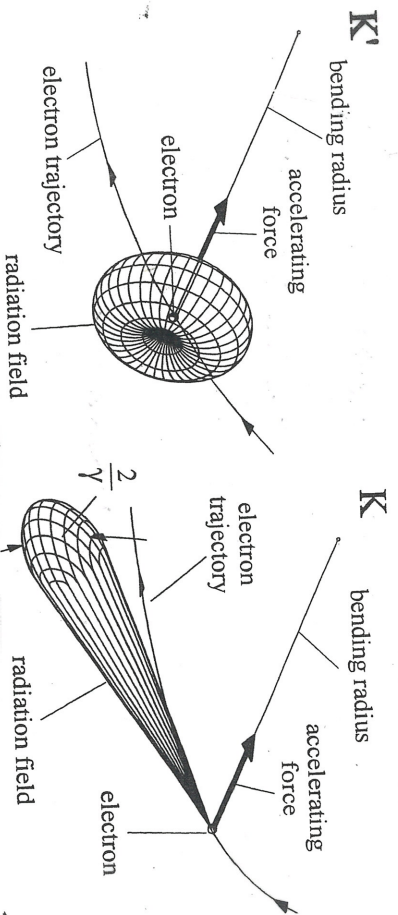
Particle frame - doughnut dipole pattern

Laboratory frame - highly Lorentz-transformed 'beam of light'

$$\gamma = 1957 \cdot E(\text{GeV}) \approx 6000 \text{ for } 3\text{GeV } e^- \text{ machines}$$

$$\frac{1}{\gamma} = \frac{1}{6000} \approx 160 \text{ radiation (1.6mm at } 10\text{m)}$$

Power density goes way up



K. White

Properties of the SR Beam

1) SR has a forward-directed angular distribution

$$\frac{dP}{d\Omega} \quad (\text{power into solid angle})$$

2) SR field has a spectral distribution

$$\frac{dP}{d\omega} \quad (\text{power into frequency } \omega)$$

3) SR field is polarized

- pure horizontal in the accelerator plane
- elliptically polarized above and below the midplane

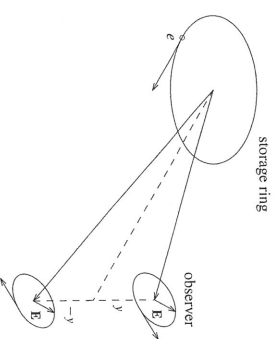
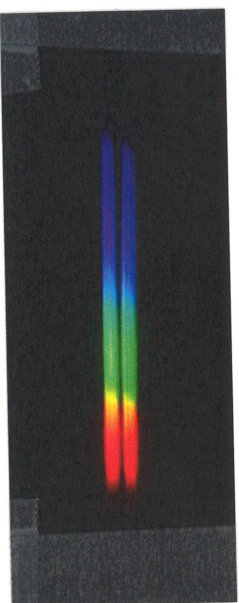
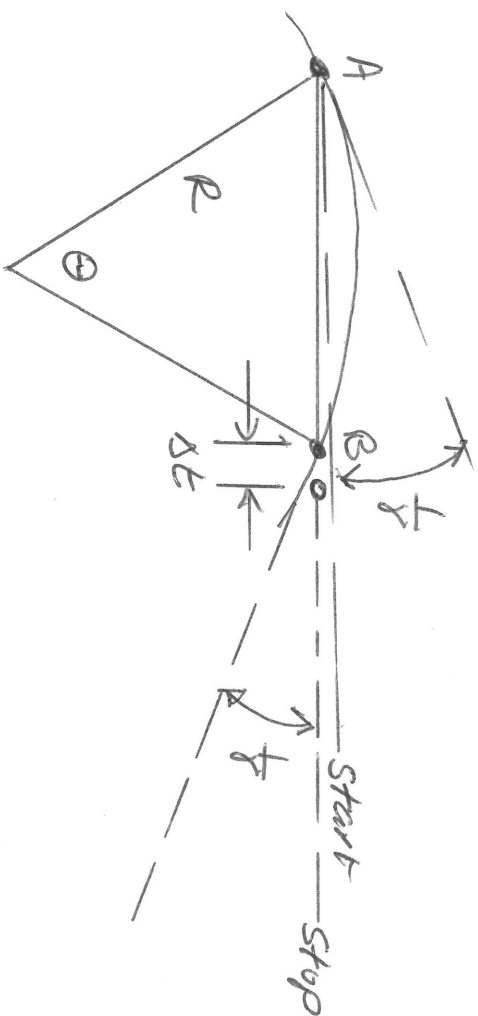


Fig. 1.6. Linear and elliptical polarization of synchrotron radiation.



Characteristic Emission Frequency



$$\Delta t = \frac{l}{c} = \frac{AB}{v} - \frac{AB}{c}$$

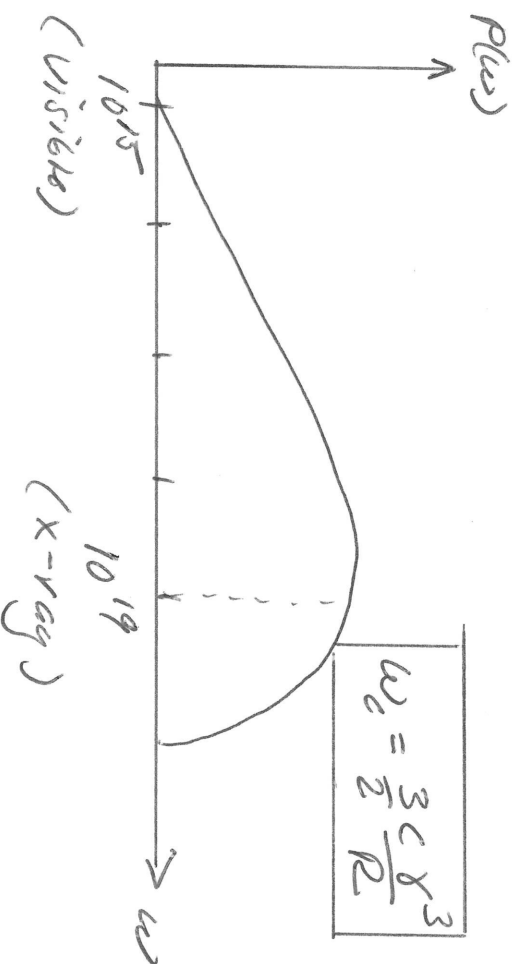
$$\Delta t = \frac{R\theta}{v} - \frac{R\sin\theta}{c}$$

$$\theta \sim \frac{l}{R} \quad \sin\theta \sim \theta = \frac{l}{R}$$

$$\Delta t \sim \frac{R}{c\gamma^3} \sim \frac{7}{3 \times 10^8 (6000)^3} \sim 10^{-19} \text{ s}$$

$$\Delta f = \frac{1}{\Delta t} = 10^{+19} \text{ Hz}$$

EM pulse from one electron
(classical treatment)

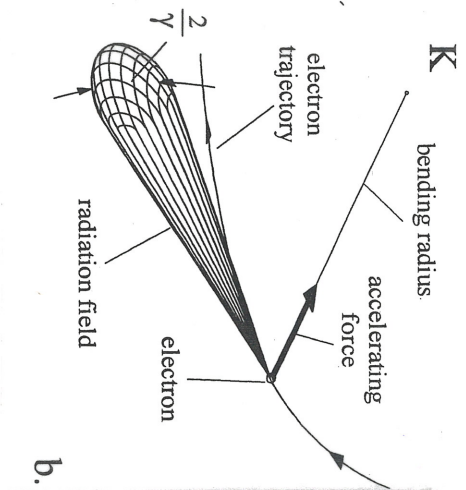


The Angular Spectral Power Density: $\frac{d^2P}{d\Omega d\omega}$

Angular

$$\frac{dP}{d\Omega}$$

$$\Theta_c \sim \frac{1}{\gamma} \sim 160 \mu\text{rad}$$



b.

Spectral

$$\frac{dP}{d\omega}$$

$$\omega_c = \frac{3}{2} \frac{c}{R} \gamma^3 \sim 10^{19}$$

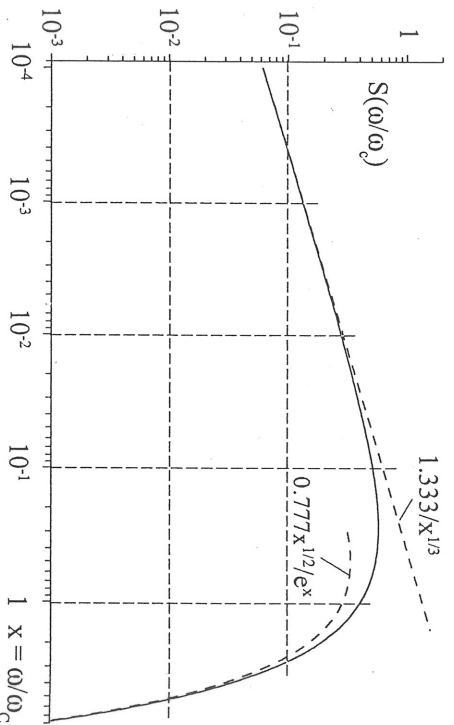
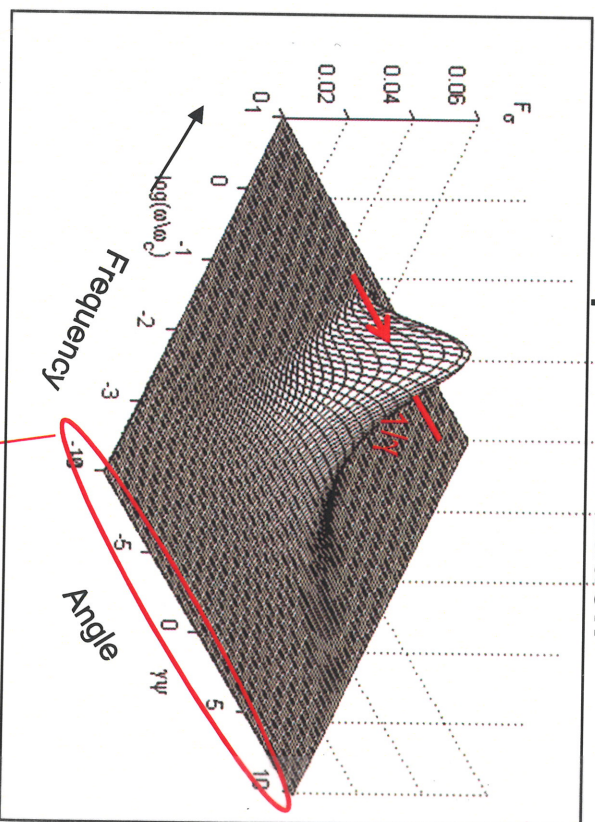
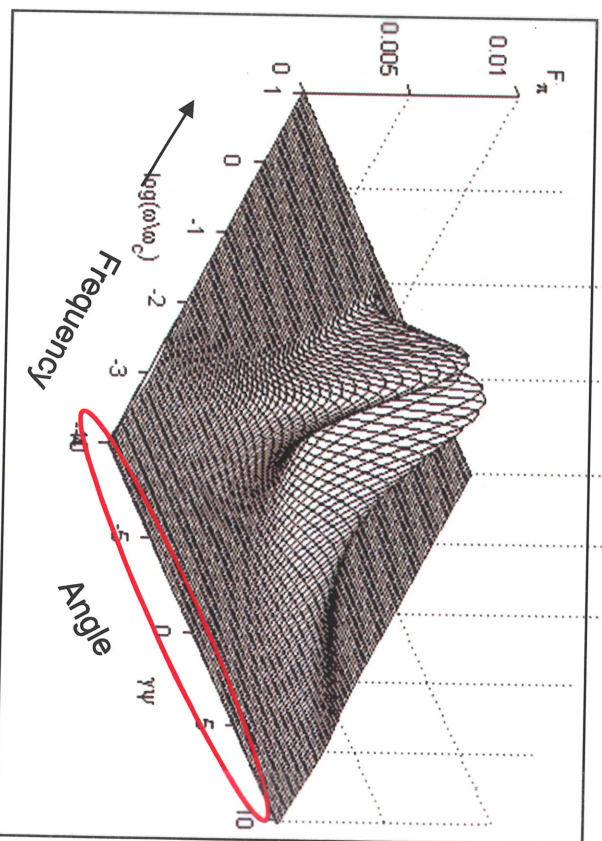


Fig. 9.5. Universal function of the synchrotron radiation spectrum $S(\omega/\omega_c)$

Horizontal Dipole Polarization

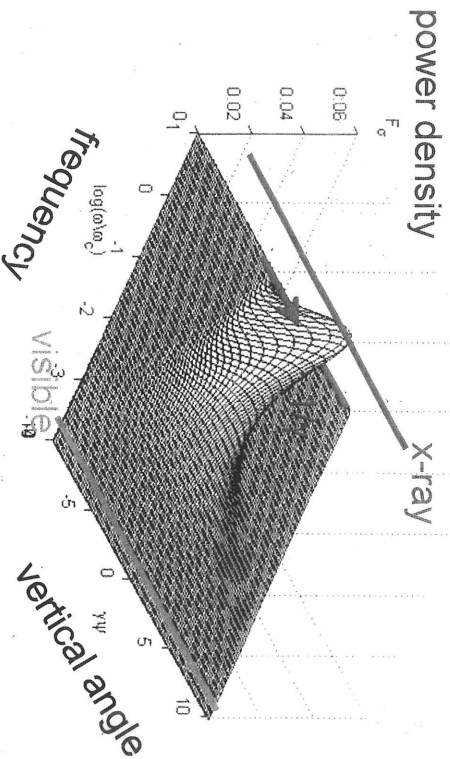


Vertical Polarization



Angular-Spectral Power Density - Schwiniger's Equations

σ-mode



$$\frac{d^2 P}{d\Omega d\omega} = \frac{P_s \gamma}{\omega_c} [F_{s\sigma}(\omega, \psi) + F_{s\pi}(\omega, \psi)]$$

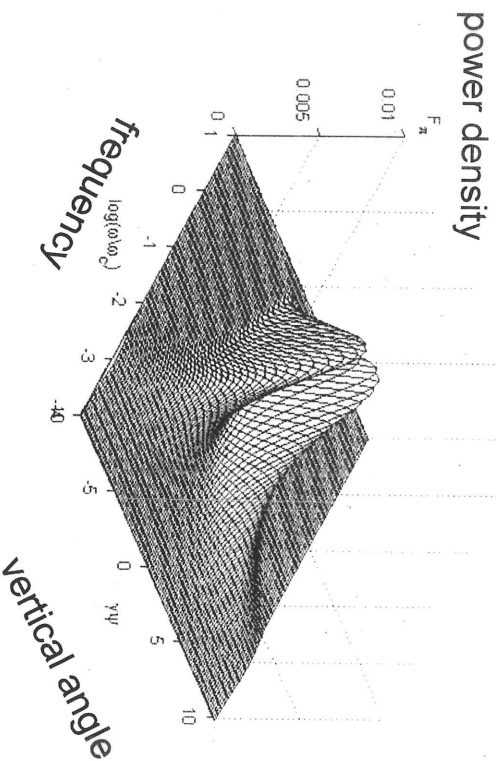
Horizontal Vertical

$$F_{s\sigma}(\omega, \psi) = \left(\frac{3}{2\pi}\right)^3 \left(\frac{\omega}{2\omega_c}\right)^2 (1 + \gamma^2 \psi^2)^2 K_{2/3}^2 \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2}\right)$$

$\frac{\omega}{\omega_c}, \gamma\psi$

modified Bessel function

π-mode



$$F_{s\pi}(\omega, \psi) = \left(\frac{3}{2\pi}\right)^3 \left(\frac{\omega}{2\omega_c}\right)^2 \gamma^2 \psi^2 (1 + \gamma^2 \psi^2) K_{1/3}^2 \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2}\right)$$

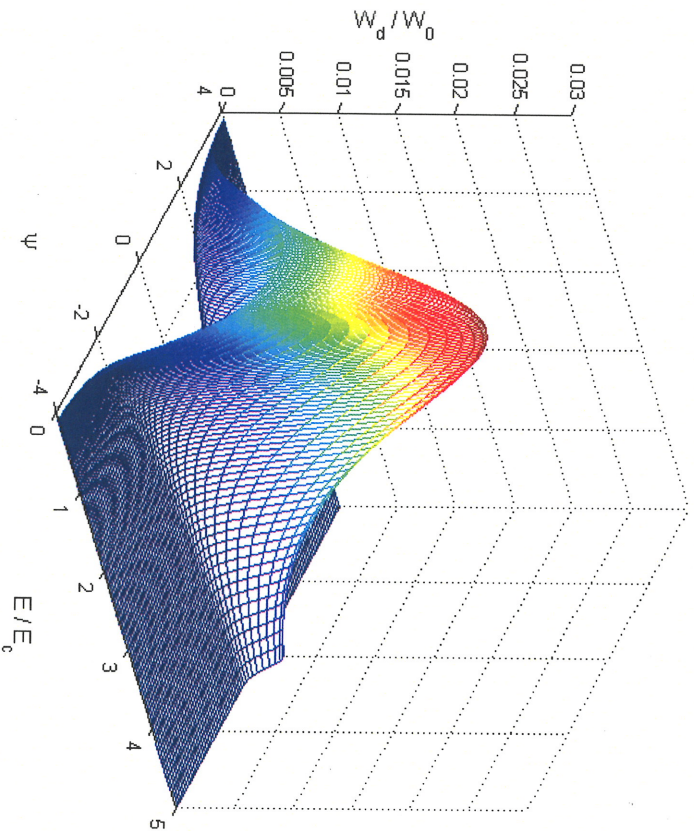
$$F_\pi = 0 \text{ at } \gamma = 0$$

Universal Curves

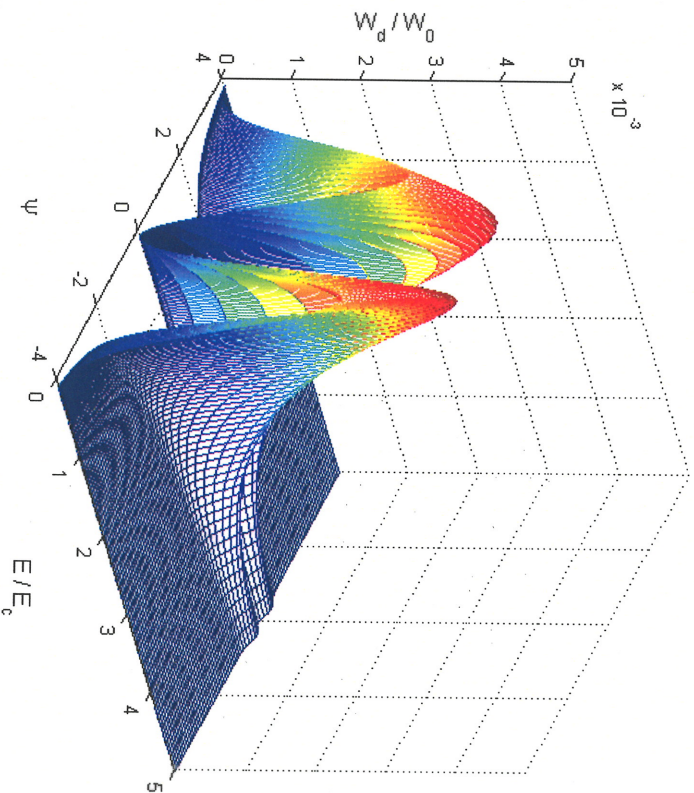
normalized to ω_c and γ

SR Beam Opening Angle

Horizontal (σ) Polarization



Vertical (π) Polarization



5.5 The angular distribution

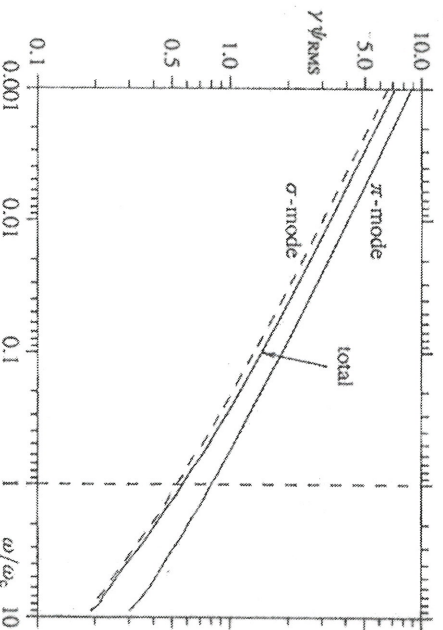


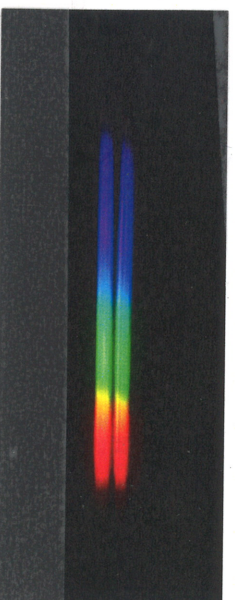
Fig. 5.8. The vertical RMS opening angle versus the frequency.

In the visible $\gamma \sim 20$

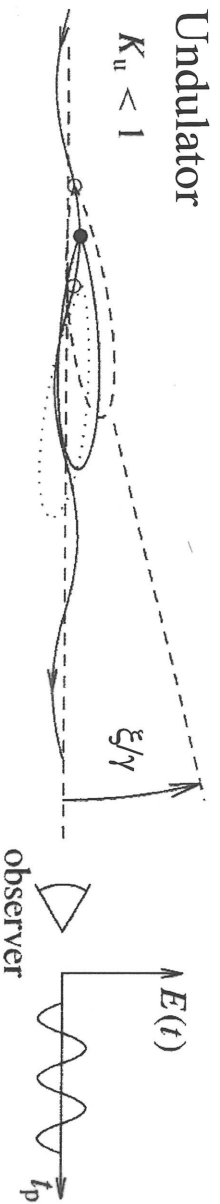
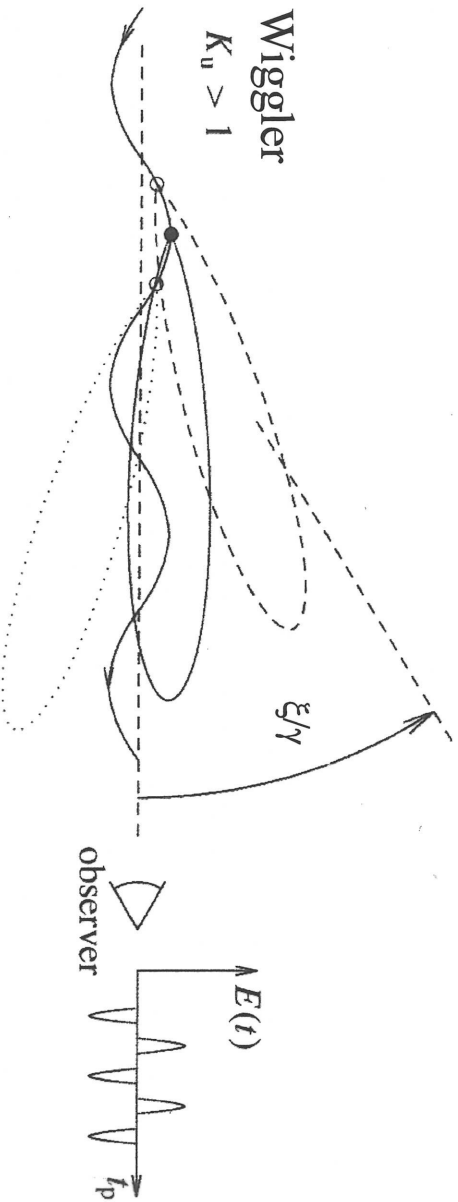
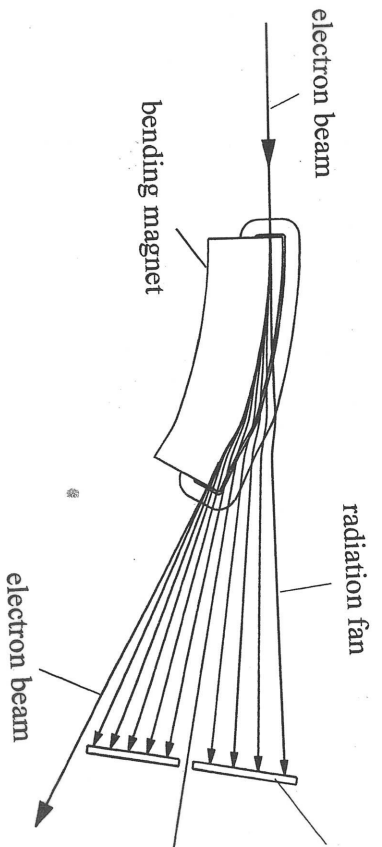
$$\gamma \sim \frac{20}{6000} = 3.3 \text{ mrad}$$

$\pm 30 \text{ mm}$ at 10 m

visible component – few mrad opening angle



SR from wigglers and Undulators



FEL: SR emission reacts back on beam $\rightarrow$ microbunching $\sigma_R \sim \lambda_{rad}$
Self-amplifies the emission process

Schwiniger's Equations - Derivation of Angular Spectral Density

$$\frac{dP}{d\omega d\Omega} = \frac{P_S}{\omega c} \left[F_\sigma(\omega, \gamma) + F_\pi(\omega, \gamma) \right] \quad \left\{ \begin{array}{l} F_{\sigma, \pi} \text{ form factors} \\ \omega \text{ frequency} \\ \gamma \text{ beam energy} \end{array} \right.$$

Power Scales as E^2

→ we need $E_\sigma(\omega, \gamma)$ and $E_\pi(\omega, \gamma)$

Start with Lienard-Wiechert Potentials

$$V(t) = \int \frac{\rho(t')}{r(t')} dV \Big|_{ret}$$

scalar potential from charge

$$\mathbf{A}(t) = \int \frac{\mathbf{j}(t')}{r(t')} dV \Big|_{ret}$$

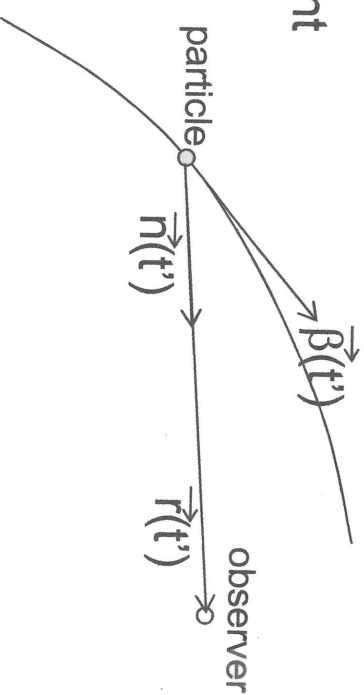
vector potential from current

$$\mathbf{E} = -\nabla V - \frac{d\mathbf{A}}{dt}$$

electric field

$$\mathbf{B} = \nabla \times \mathbf{A}$$

magnetic field



Schuringer's Equations

$$\mathbf{E} = -\nabla V - \frac{d\mathbf{A}}{dt}$$

$$\mathbf{B} = \nabla \times \mathbf{A}$$

After a Jacksonian derivation

$$\mathbf{E}(t) = \frac{1}{4\pi\epsilon_0 c} \left\{ \frac{\left[\mathbf{n} \times [(\mathbf{n} - \boldsymbol{\beta})] \times \dot{\boldsymbol{\beta}} \right]}{r(1 - \mathbf{n} \cdot \boldsymbol{\beta})^3} \right\}_{\text{ret}}$$
$$\mathbf{B}(t) = \frac{\mathbf{n} \times \mathbf{E}}{c} \Big|_{\text{ret}}$$

Poynting Vector (power flux)

'difficulty evaluating the above equations'
'advantageous to calculate Fourier transforms'

Fourier Transform Fields + Small Angle Approximations

$$\tilde{E}_x(\omega) = \frac{-\sqrt{3}e\gamma}{(2\pi)^{3/2} \epsilon_0 c r_c} \cdot \frac{\omega}{2\omega_c} \cdot (1 + \gamma^2 \psi^2) K_{2/3} \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2} \right)$$

$$\tilde{E}_y(\omega) = \frac{i\sqrt{3}e\gamma}{(2\pi)^{3/2} \epsilon_0 c r_c} \cdot \frac{\omega}{2\omega_c} \cdot \gamma\psi (1 + \gamma^2 \psi^2) K_{1/3} \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2} \right)$$

there are two polarizations

$$\frac{d^2 P}{d\Omega d\omega} = \frac{d^2 P_\sigma}{d\Omega d\omega} + \frac{d^2 P_\pi}{d\Omega d\omega} = \text{Const} \cdot \left(|\tilde{E}_x(\omega)|^2 + |\tilde{E}_y(\omega)|^2 \right).$$

$$F_{s\sigma}(\omega, \psi) = \left(\frac{3}{2\pi} \right)^3 \left(\frac{\omega}{2\omega_c} \right)^2 (1 + \gamma^2 \psi^2)^2 K_{2/3}^2 \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2} \right).$$

$$F_{s\pi}(\omega, \psi) = \left(\frac{3}{2\pi} \right)^3 \left(\frac{\omega}{2\omega_c} \right)^2 \gamma^2 \psi^2 (1 + \gamma^2 \psi^2) K_{1/3}^2 \left(\frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2} \right).$$

Radiated Power and Energy Loss per Turn

Previously $P_{rad} \propto E^2 B^2$ or $\frac{E^4}{R^2}$ $E = \text{Beam Energy}$
 $B = \text{Dipole Field}$
 $R = \text{Bend radius}$

P_{rad} = instantaneous power radiated by one electron

Energy loss per revolution

$$U_0 = \oint P_{rad} dt = \frac{1}{c} \oint P_{rad} ds \propto \frac{E^2}{c} \oint B^2(s) ds$$

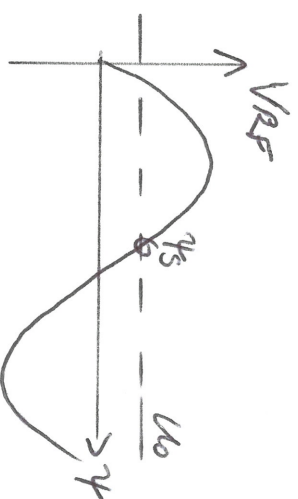
Bend radius $\frac{1}{R} = \frac{gB}{8mc^2}$

$$U_0 \propto E^2 \oint \frac{ds}{R^2(s)}$$

↑ Synchrotron integral I_2 ($I_1, -I_5$)

Get relativistic factors correct

$$U_0 = \frac{e^2 \gamma^4}{3 \epsilon_0 R} = \frac{88.5 E^4 (\text{GeV})}{R(m)} \text{ keV}$$



$$U_0 \sim \frac{100 \cdot 100}{10} = \frac{10^4}{10} = 10^3 \text{ keV} \\ (1 \text{ MeV/turn})$$

SR Power and Energy Loss Scaling

Power

$$P_{rad} \propto \frac{\gamma^4}{R^2}$$

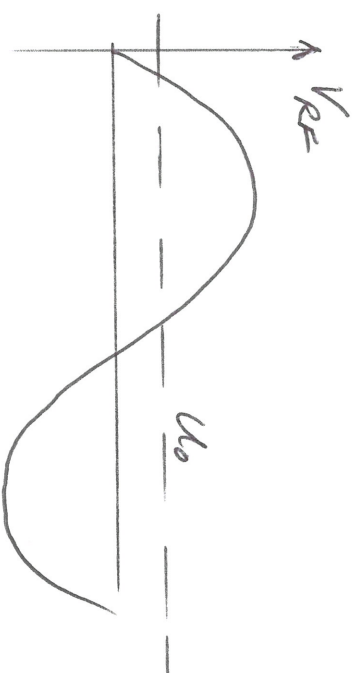
Joules
sec

Energy Loss
Turn

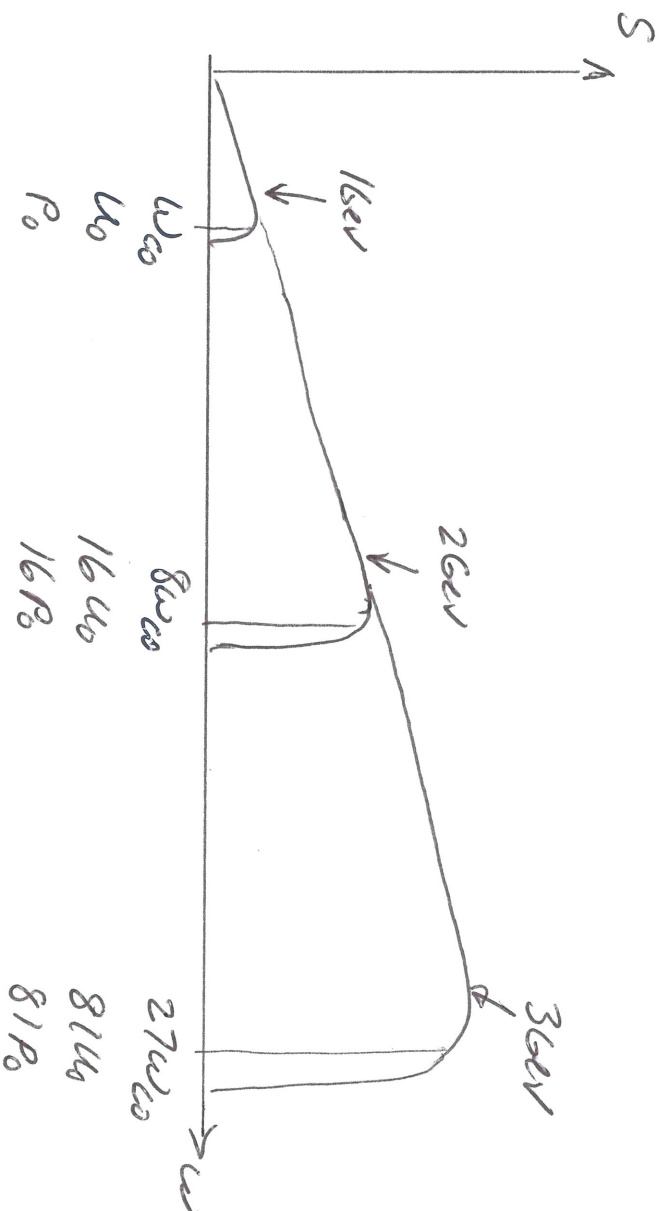
$$U_0 \propto \frac{\gamma^4}{R}$$

Joule (eV)

$$\left(P = \frac{U_0}{T_0} \text{ where } T_0 = \frac{2\pi R}{c} \right)$$



Critical Energy $\omega_c = \frac{3}{2} \frac{c \gamma^3}{R}$



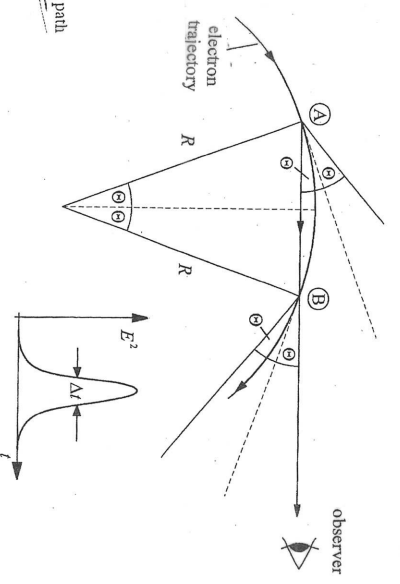
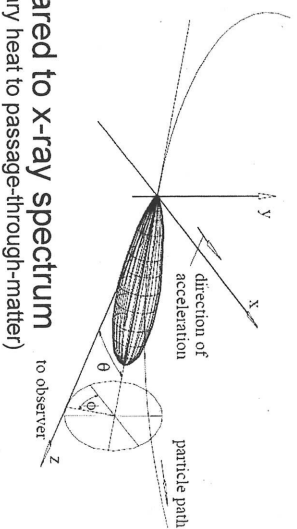
Harder spectrum
more power

SR Review

1) SR spectrum extends into hard X-ray

2) Radiation forward-directed

infrared to x-ray spectrum
(ordinary heat to passage-through-matter)



3) Angular-Spectral power density $\frac{d^2 P}{d\Omega d\nu}$
(Schwinger's Equations)

4) Horizontal/Vertical Polarization components (or left/right circular)

5) Critical photon energy scales as $\frac{\gamma^3}{R}$

6) Instantaneous radiation power and energy loss/turn scale as γ^4